

A rate-state model for aftershocks triggered by dislocation on a rectangular fault: a review and new insights

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Abstract

We compute the static displacement, stress, strain and the Coulomb failure stress produced in an elastic medium by a finite size rectangular fault after its dislocation with uniform stress-drop but a non uniform dislocation on the source. The time-dependent rate of triggered earthquakes is estimated by a rate-state model applied to a uniformly distributed population of faults whose equilibrium is perturbed by a stress change caused only by the first dislocation. The rate of triggered events in our simulations is exponentially proportional to the shear stress change, but the time at which the maximum rate begins to decrease is variable from fractions of hour for positive stress changes of the order of some MPa, up to more than a year for smaller stress changes. As a consequence, the final number of triggered events is proportional to the shear stress change. The model predicts that the total number of events triggered on a plane containing the fault is proportional to the $2/3$ power of the seismic moment. Indeed, the total number of aftershocks produced on the fault plane scales in magnitude, M , as 10^M . Including the negative contribution of the stress-drop inside the source, we observe that the number of events inhibited on the fault is, at long term, nearly identical to the number of those induced outside, representing a sort of conservative natural rule. Considering its behavior in time, our model does not completely match the popular Omori Law; in fact it has been shown that the seismicity induced closely to the fault edges is intense but of short duration, while that expected at large distances (up to some tens times the fault dimensions) exhibits a much slower decay.

Key words *aftershocks – earthquake triggering – fault parameters – rate-and-state – Omori Law*

1. Introduction

Studying the space-time interaction of earthquakes is important not only for the comprehension of earthquakes, but also for its possible application to the assessment of seismic hazard.

In this paper we revisit the physical modeling of the interaction between seismic events, seeking a relationship between the source parameters of an earthquake and the rate of all other earth-

quakes that follow it. The results obtained from the simulations performed with the physical model proposed in this study are not totally innovative; however, they are rarely taken into account by the numerous studies concerning the fault interaction and earthquake triggering.

A purely statistical approach to this phenomenon exists with the name of «epidemic model» (Ogata, 1988, 1998; Console *et al.*, 2003). In the epidemic model each earthquake source is supposed capable of increasing the probability of new earthquakes according to a spatio-temporal kernel representing the probability distribution, which is decreasing from the causative source. The epidemic model is applicable not only to the aftershock phenomenon but also to the foreshock. However, although it provides an accurate statistical representation of earthquake interaction, the physical interpretation is neglected.

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In order to formulate a more physical description of clustered seismicity, we consider that seismic events modify the stress field around the causative fault. Our main goal in this work is to observe the spatio-temporal variation of the seismicity rate after an event introducing in the algorithm a rate-state model approach derived by Ruina (1983) and Dieterich (1986, 1992, 1994). In particular, we seek a physically based relationship between the magnitude of the causative earthquake and the number of its aftershocks. Some recent studies disclosed that sudden stress variations, even of small magnitude, can produce large variations of the seismicity rate. It is recognized as a phenomenon of triggering on faults that are already loaded by tectonic stress and in close-to-failure condition. The seismicity rate increases in general where the stress change (called Coulomb Stress Change) is positive, according to the Coulomb model (Mendoza and Hartzell, 1988; Boatwright and Cocco, 1996; Stein *et al.*, 1997; Gomberg *et al.*, 1998, 2000; Harris, 1998; Stein, 1999; Toda and Stein, 2000; King and Cocco, 2001; Kilb *et al.*, 2002; Belardinelli *et al.*, 2003; Nostro *et al.*, 2005). These studies were able to give a physical interpretation for earthquake interaction observed in specific real cases, but were not suitable for the general application of the model in a predictive way that can be tested for example using the information contained in a seismic catalog. In particular, the Coulomb stress change criterion does not model the hyperbolic time decay of the aftershock rate after a mainshock, known as the Omori Law. In this simple model all the impending earthquakes would be clockadvanced by a constant time that depends on the stress change and on the loading rate but does not depend on the physical properties of the fault (Gomberg *et al.*, 1998). The experience shows that the Earth does not rather behave in this way. The most popular empirical description of the phenomenon is the Omori Law, stating that the aftershock rate decays in time as t^{-1} . Among the various theories modeling the Omori Law, we have taken into consideration only the one that assumes the rate-and-state dependent constitutive model of faults. We have considered only the static stress changes created by an earthquake. We have not considered, even without rejecting them in principle, other hypotheses that predict a time

variation in stress, like viscoelastic relaxation or the diffusion of fluids in the crust. The rate-state model for earthquake nucleation seems capable of substantially explaining all the phenomenology. In this model a promoted failure does not occur instantaneously at the exceeding of a threshold stress as for the Coulomb-Amonton criterion, but follows a more complex non-linear time history with different phases. This time history depends on the physical properties of the fault too. Gomberg *et al.* (2005) have shown that the rate-state Dieterich model may be derived from simple general ideas. The application of the rate-and-state model to the evolution of seismicity is not new; Toda *et al.* (1998) used this model to predict the spatio-temporal seismicity rate after the Kobe earthquake. Other examples of application of the same model can also be found in the literature (Toda and Stein, 2000, 2002, 2003). Following the rate-state model it has recently been possible to simulate seismicity quite realistically, accounting for the rate of events triggered by stress changes even at large distance from the inducing earthquake (Ziv, 2003; Ziv and Rubin, 2003).

Basing our work on the classical theory of the elasticity, we start from the formulation of the stress released by a rectangular fault with uniform stress-drop; then we apply the rate-state model (Dieterich, 1994) obtaining the complete time and space distribution of the seismicity induced by a given earthquake. It is important to remark that for present purposes we neglect the interaction between subsequent events, supposing that only the first event perturbs the stress field. We also neglect the effect of the free surface of the Earth. In this study we focus on the behavior of induced seismicity in space (on the fault plane) and time and its relationship with the fault parameters. The results of simulations are analyzed to find the scaling relationships existing among the free parameters of the model and the expected seismicity in a way that will allow the validation of the model by real observations.

2. Elastic model for a rectangular fault

The theoretical approach of this work has been made as simple as possible, still preserv-

ing the capability of modeling the stress transfer from a seismic source to another.

Suppose that a fault is embedded in a homogeneous, isotropic, elastic medium where the stress tensor σ is uniform. At a particular moment, the fault slips generating an earthquake, and suppose that the earthquake fully releases the component of the traction parallel to the slip vector across the fault. We call $\Delta\sigma$ (the stress-drop) the uniform negative change of traction on the fault. It is well known that the slip distribution on a fault in similar conditions is not uniform. The slip distribution will satisfy the theory of the elasticity at the new equilibrium, and it will be zero on the edges of the fault. It is also well known that the shear stress on the fault plane outside the edges of the fault increases significantly.

The analytical solution for the slip distribution on a fault at equilibrium does not exist except that for very simple geometries, such as a rectangular fault of constant width and infinite length (Knopoff, 1957) and for a circular fault (Keilis-Borok, 1959; Udias, 1999). As mentioned by Kostrov and Das (1998), the analytical solution is not known even for a geometry as simple as a rectangular shape of width W and length L . In the literature it is possible to find the computation of the stress field generated by an ideal rectangular fault with uniform slip in an infinite homogeneous space or halfspace (Chinnery, 1961; Iwasaki and Sato, 1979; Okada, 1985, 1992). However, as mentioned above, this is not a realistic model since it implies the compenetrability of the medium at the edges of the slip region. For this reason, here we make use of a non uniform slip distribution that is approximatively compatible with a uniform stress-drop on the fault.

To calculate the static displacement, stress and strain, we start from the Somigliana Tensor and we make use of the double couple point source elastic theory. We compute the modified Coulomb failure stress from the relation

$$\sigma_c = T_s + \mu' T_n \quad (2.1)$$

where T_s and T_n are the traction parallel to the slip vector and the traction normal to the fault plane (negative for compression) respectively; μ' is the effective coefficient of friction incorporating the pore fluid presson that is commonly assumed to

be proportional to normal stress change (although this assumption is not general, see Cocco and Rice, 2002). The effective coefficient of friction in coseismic stress change studies varied from 0.0 to 0.75, with an average value of $\mu'=0.4$ often used. This is the value we used in our simulations.

Our model of an extended source is drawn heuristically from that of a rectangular fault infinitely extending in one direction and that of a circular fault, and it is characterized by a slip distribution as

$$\Delta u = \frac{\Delta\sigma}{\mu} \sqrt{\frac{\left[\left(\frac{W}{2}\right)^2 - x^2\right]\left[\left(\frac{L}{2}\right)^2 - y^2\right]}{\frac{LW}{4}}} \quad (2.2)$$

defined for $-W/2 \leq x \leq W/2$ and $-L/2 \leq y \leq L/2$, where x, y are the coordinates on the fault plane whose origin is coincident with the center of the rectangle, $\Delta\sigma$ is the stress-drop and μ the shear modulus (rigidity) of the elastic medium. Equation (2.2) satisfies the condition of zero slip on the edges of the rectangle. The slip distribution is of a form very similar to that obtained numerically by Kostrov and Das (1998) with the finite differences method. Its analytical form is similar to the solution for a circular fault. Moreover, for a fault of square shape ($W=L$) it achieves a seismic moment consistent with those obtained analytically for two circular faults, respectively inscribed and circumscribed to the square. Bonafede and Neri (2000) have shown that, when imposing a unidirectional traction release in the strike or dip direction, a minor component of slip is present, over the fault plane, even in the direction perpendicular to the released traction. We have neglected this component.

For our numerical applications we discretize the continuous slip distribution of eq. (2.2) with a set of point sources densely and uniformly distributed on the fault. The more numerous the point sources are, the more accurate the approximation will be.

In our numerical simulations, we arbitrarily, but without loss of generality, suppose that the fault is on the (x_1, x_3) plane. Let us note that some fault parameters are connected among

each other; in particular, the scalar moment M_0 is defined through the relation

$$\begin{aligned} M_0 &= \mu \overline{\Delta u} WL = \frac{\pi^2}{16} \mu \Delta u_{\max} WL = \\ &= \frac{\pi^2}{32} \Delta \sigma (WL)^{3/2} \end{aligned} \quad (2.3)$$

where $\overline{\Delta u}$ is the average slip on the fault and the maximum slip, Δu_{\max} , is defined as

$$\Delta u_{\max} = \frac{\Delta \sigma}{\mu} \frac{\sqrt{LW}}{2}. \quad (2.4)$$

We also assume that the magnitude M is proportional to the logarithm of the seismic moment M_0 (measured in Nm) according to the Kanamori and Anderson (1975) relation

$$\log M_0 = 9.1 + 1.5M. \quad (2.5)$$

The computer code allows the user to chose among several options:

- The focal mechanism characterizing the causative fault and the induced events constrained by the strike ϕ , the dip δ and the rake λ angles.
- Dimensions and stress-drop, seismic moment, or magnitude of the causative fault.
- Type of output: displacement (3 components), stress or strain (6 independent components), traction parallel to the slip vector and

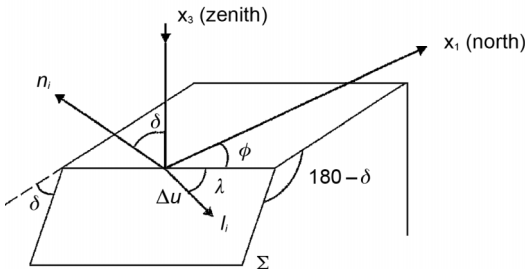


Fig. 1. Conventions adopted for the reference system and for the parameters of the orientation of a fault: the strike ϕ , dip δ and rake λ angles and the unit vectors \mathbf{n} and \mathbf{l} . The x_3 axis is positive toward the Nadir direction.

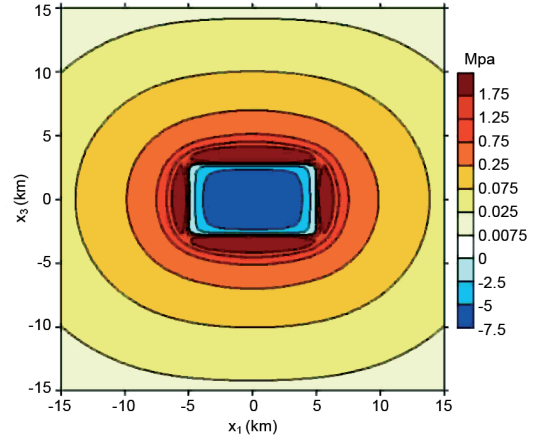


Fig. 2. Shear stress, σ_{12} , parallel to the slip direction for a strikeslip fault projected on the fault plane (x_1, x_3 , $x_2=0.2$ km).

Table I. Source parameters used in numerical applications.

Parameter	Value
Strike, dip and rake	0, 90, 0
Dimension of the fault	10 km \times 6 km
Spacing on the fault	0.1 km
Dimension of the grid	30 km \times 30 km
Spacing on the grid	0.5 km
Magnitude	6

perpendicular to the fault plane (2 component) and Coulomb failure stress (1 component).

The output values are given on a rectangular grid of points, whose orientation (horizontal, or vertical parallel to one of the coordinate axes), dimensions and spacing are also selected by the user. See fig. 1 for the conventions we adopted for the reference system and the angles.

For sake of simplicity, in this work we make the assumption that all the aftershocks have the same focal mechanism of the causative earthquake.

Figure 2 shows the component of the stress field σ_{12} parallel to the slip calculated through the parameters listed in table I. Note that this

figure shows the stress change and not its absolute value.

Figure 2 clearly shows that the shear stress σ_{12} is maximum near the edges of the fault, in particular along the slip direction x_1 , and decreases with increasing the distance from the fault. Inside the fault the stress change is negative and approximately equal to the imposed value of $-\Delta\sigma$.

3. The rate-state model

As already stated, we assume that the time-dependent behavior of the seismicity triggered by a shear stress change in a population of faults is described by the rate-state model introduced by Dieterich (1992, 1994) where the constitutive model proposed by Ruina (1983) and Dieterich (1986) is used. According to this model the rate R of earthquakes due to a $\Delta\tau$ shear stress perturbation is given by

$$R = \frac{r \frac{\dot{\tau}}{\dot{\tau}_r}}{\left[\frac{\dot{\tau}}{\dot{\tau}_r} \exp\left(\frac{-\Delta\tau}{A\sigma}\right) - 1 \right] \exp\left(\frac{-t}{t_a}\right) + 1} \quad \dot{\tau} \neq 0 \quad (3.1)$$

whose various parameters are defined in table II. Note that eq. (3.1) does not account for the magnitude of the earthquakes, for the distribution of which we assume the classical Gutenberg-Richter frequency-magnitude relationship. However, both R and r refer to the number of earthquakes in a specific magnitude range.

We make the assumption $\dot{\tau} = \dot{\tau}_r$. The meaning of this assumption is that the primary earth-

Table II. Constitutive parameters used in Dieterich's relation.

Symbol	Description
r	reference seismicity rate
$\dot{\tau}$	shear stressing rate
$\dot{\tau}_r$	reference stressing rate
$\Delta\tau$	earthquake shear stress change
A	fault constitutive parameter
σ	normal stress
t	time
t_a	characteristic time for seismicity to return to the steady state

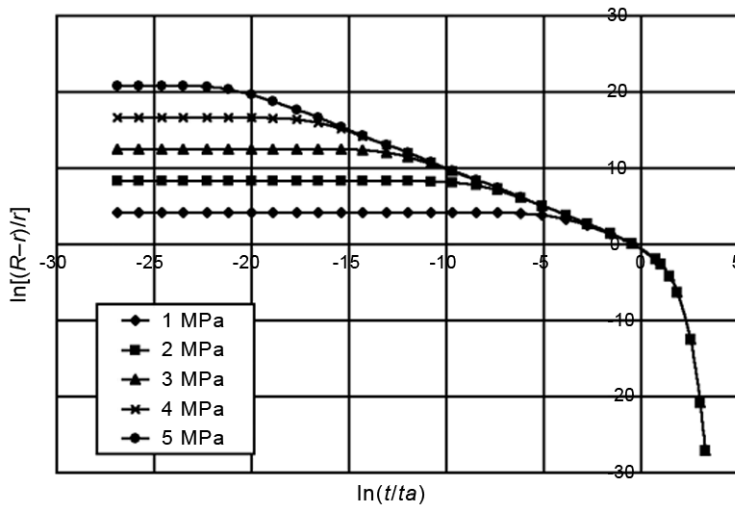


Fig. 3. Distribution of the triggered seismicity against time for different values of induced shear stress. The values of the parameters used are: $r=2$ events/yr, $A\sigma=0.24$ MPa and $\dot{\tau}=5$ KPa/yr.

quake does not change the long-term stress rate due to the tectonic driving forces in the environment.

It is easily understandable from (3.1) that the integral of R over infinite time diverges. This is due to the fact that the limit of R for $t \rightarrow \infty$ is the background rate r . We prefer to deal with the net triggered seismicity rate $R' = R - r$. The plot of R' versus time for different values of $\Delta\tau$ is shown in fig. 3 (see also Dieterich, 1994). It shows that the initial value of the triggered seismicity rate, immediately after the triggering event, is proportional to the exponential of the stress variation but the time period over which $R(t)$ is maximum and relatively constant decreases exponentially by the same factor for increasing shear stress change. This time period, in substantial agreement with Dieterich (1994), can be obtained computing the intersection between $R'(t=0)$ and the straight line corresponding to the limit of R' when $\Delta\tau \rightarrow \infty$. The difference with Dieterich (1994) is that we have preferred to make use of R' , rather than R , and we ignore the case $\tau = 0$ which we judge not realistic. Such time period ranges from few hours

for a variation $\Delta\tau = 5$ MPa to nearly one year for $\Delta\tau = 1$ MPa. Figure 3 finally shows that the total duration of the triggered seismicity is the same but the shape of the decay is different for different $\Delta\tau$.

The integration of R' over infinite time leads to the not completely trivial result that the total number of triggered events in an area of constant stress change is proportional to $\Delta\tau$

$$\int_0^\infty R'(t) dt = \frac{rt_a}{A\sigma} \Delta\tau = \frac{r}{\dot{\epsilon}} \Delta\tau. \quad (3.2)$$

This relationship, predicted by the friction model, had been observed for «Landers Aftershocks» by Gross and Kisslinger (1997).

Consider the integral over time of the triggered seismicity rate R' , once for a positive $\Delta\tau$, and then for a negative $\Delta\tau$ of the same size, and we sum the two functions. We obtain a quantity that initially has a positive increase and then tends back to zero (fig. 4). This result can be interpreted as a sort of conservative law of the seismicity, by saying that the expected total number of aftershocks generated in a place close to a fault can be balanced, at long range, by the total

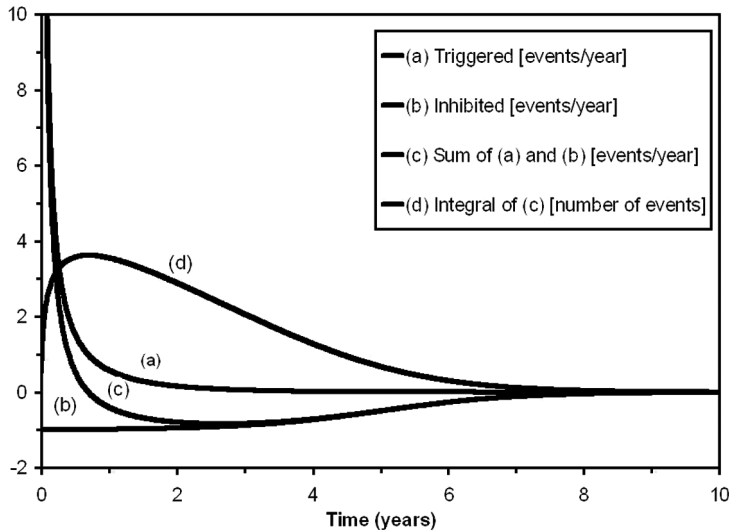


Fig. 4. Distribution of the number of expected events in four different cases: (a) for a positive value of earthquake shear stress equal to 5 MPa, (b) for a negative value of it equal to -5 MPa, (c) the sum of (a) and (b) and (d) the integral of (c) over time. The values of r and t_a are fixed at 1.

number of those events that are inhibited by the stress-drop in a place internal to the fault if the stress-drops are equal in absolute value. It seems to us that this circumstance has been ignored, so far, in all previous studies on this subject.

For numerical applications it is necessary to define the value of the various parameters appearing in eq. (3.1). Table III contains a list of values inferred from several geophysical considerations and used in our simulations. The reference seismicity rate, r , is a quantity that should be known experimentally for a given area, and it is related to the average strain rate $\dot{\epsilon}$ of that area. We use for r a uniform value of four events of magnitude exceeding 3.0 per year per 1000 km², and for $\dot{\epsilon}$ a value of 5 KPa per year. These are simplifications of the reality where r and $\dot{\epsilon}$ are not geographically uniform. These parameters seem however reasonable for an area of moderate seismicity. The reason r is expressed as a number of events per unit area and time is given below. The parameter A of the constitutive law has a value ranging from 0.005 to 0.015 obtained from laboratory experiments (Dieterich, 1995). Generally, simulations that most nearly resemble earthquakes in nature were obtained with rather small values of A , near to $A \approx 0.001$. The value of $A\sigma$ has been evaluated for different sequences of earthquakes by several authors (Stein, 1999; Belardinelli *et al.*, 2003). The parameter σ is the normal stress on the fault (of the order of some tens of MPa). The characteristic time of aftershocks, t_a , depending on the other parameters through the relation $t_a = (A\sigma)/\dot{\epsilon}$, in this context is of the

order of several tens of years. This value is essential to model a long term triggered seismicity, because the effect of the stress change disappears completely after a time of the order of the double of t_a .

4. Application to a rectangular fault

Dieterich (1994) discussed in detail the case of a finite circular dislocation with uniform stress-drop. In our simulations we have taken into consideration a vertical, rectangular strike-slip fault embedded in an infinite, homogeneous, isotropic, elastic medium. The values assumed for the fault parameters are again those reported in table I. Assuming a constant stress-drop $A\sigma$, the scalar seismic moment M_0 is proportional to the quantity $(WL)^{3/2}$, that can be considered a sort of equivalent source volume. Scaling W and L by a given factor, M_0 scales accordingly, and also the shape of the stress distribution in space scales by the same factor. Assuming that the background seismicity rate is given as the number of events per unit volume and time, and integrating over the space, we find that the number of events triggered by the earthquake is proportional to the source volume, *i.e.* to the seismic moment M_0 , or to $10^{3/2M}$, where M is the magnitude. This study analyzes only those aftershocks occurring on the fault plane neglecting off-fault seismicity and assuming that this is the plane of weakness optimally oriented with respect to the tectonic loading stress. The modeling of the space-time distribution of seismicity triggered by an earthquake is straightforward, introducing the expression for the Coulomb failure stress σ_c in eq. (3.1) for $\Delta\tau$ considering that on the fault plane $\sigma_c \equiv \Delta\tau$. In fact, on the fault plane, the normal stress changes in a homogeneous and isotropic medium are likely to vanish.

In this case, the number of triggered events would be proportional to the fault area of the triggering event and then to $M_0^{2/3}$, or to 10^M . If this scaling is true, one earthquake of magnitude M would have a number of aftershocks ten times larger than one earthquake of one magnitude unit smaller. Assuming, then, that the b -value of the Gutenberg-Richter frequency-mag-

Table III. Values of the parameters used in the numerical applications of the model for induced seismicity.

Parameter	Value
r	4 events/(yr·1000 km ²)
$\dot{\epsilon}$	5 KPa
$\dot{\epsilon}_r$	5 KPa
A	0.008
σ	30 MPa
t_a	48 yr

nitude relation is equal to 1, it would follow that the earthquakes of magnitude M , all together, produce the same number of aftershocks as the earthquakes of one magnitude unit smaller. This hypothesis is supported by the statistical analysis of real observations (Utsu, 1969; Yamanaka and Shimazaki, 1990; Felzer *et al.*, 2001; Helmstetter, 2003; Helmstetter *et al.*, 2005) and will be used in the following.

A first group of simulations concerning the seismicity triggered by a fault as a whole have been performed to study the variation of the total number of triggered events (in space and time), N_{tot} , versus the main free source parameters: the stress-drop $\Delta\sigma$, the linear dimensions L and W and the magnitude M .

Tables IV, V and VI show the results obtained assigning a fixed value for one of the parameters at a time, and letting the others change according to eqs. (2.3) and (2.4). These simula-

Table IV. Relation between fault parameters and aftershock production. The value of $\Delta\sigma$ is fixed at 4 MPa; the grid used in this computation has an area of 10^4 km², the spacing on it is 0.5 km. The spacing on the fault is 0.05 km.

$L \cdot W$ [km ²]	M	M_0 [Nm]	N_{tot}
4	4.6	$9.9 \cdot 10^{15}$	13.4
25	5.4	$1.5 \cdot 10^{17}$	78.2
100	6.0	$1.2 \cdot 10^{18}$	263.2
225	6.4	$4.2 \cdot 10^{18}$	532.8

Table V. Relation between fault parameters and aftershock production. The area of the fault is fixed at the value of 100 km², the spacing is of 0.05 km; the grid used in the computation has an area of 17×17 km², the spacing on it is of 0.5 km.

$\Delta\sigma$ [MPa]	M	M_0 [Nm]	N_{tot}
1	5.6	$3.1 \cdot 10^{17}$	51.8
2	5.4	$6.2 \cdot 10^{17}$	103.4
3	5.9	$9.3 \cdot 10^{17}$	155.2
4	6.0	$1.2 \cdot 10^{18}$	206.8
5	6.1	$1.6 \cdot 10^{18}$	258.6

Table VI. Relation between fault parameters and aftershock production. The value of M is fixed at 6.0; the output grid is of 20 km \times 20 km, the spacing is 0.5 km. The spacing on the fault is 0.05 km.

$L \cdot W$ [km ²]	$\Delta\sigma$ [MPa]	N_{tot}
4	510.7	Over-flow
25	32.7	597.6
100	4.1	223.2
225	1.2	107.8

tions have been carried out by summation of the single contributions of the elementary cells constituting the area surrounding the fault. Only the seismicity triggered on the fault plane, but outside the fault edges, has been considered. Looking at table IV one can note that, under our hypotheses, for a constant stress-drop, the total number of triggered events increases proportionally to the fault area LW . Theoretically it should be exactly proportional to the area of the source, but in the calculations there are numerical imprecisions due to the grid size approximations. These numerical simulations show, also, that the number of aftershocks is proportional to the $2/3$ power of the scalar seismic moment or to the exponential of the magnitude (10^M), in agreement with the above theoretical arguments. Table V clearly shows the proportionality existing between the stress-drop $\Delta\sigma$ and N_{tot} , having fixed the linear dimensions of the fault. This is a simple consequence of the proportionality between the stress-drop, $\Delta\sigma$, and the shear stress change, $\Delta\tau$, at any point of the space around the source; furthermore, it is an example of the proportionality expressed in eq. (3.2) between $\Delta\tau$ and the total number of triggered events, integrated on the whole space. The results reported in table VI give credit to the hypothesis that, for a constant magnitude, the earthquakes of smaller linear dimensions generate a larger number of triggered events. Indeed, a point source would yield the largest quantity of triggered events for the same magnitude.

Let us mention, lastly, the results of a simple numerical test. We have theoretically shown in the previous section that the total number of

events triggered by a stress change $\Delta\tau$ per unit area is, after a time much larger than the characteristic time t_a , identical to the number of events inhibited by a negative shear stress change $-\Delta\tau$ on the same area. The numerical simulations on an extended source, with parameter still shown in table I, reflect this idea in a more general way. In fact, if we take the total number of aftershocks produced in a time interval of 100 years by a fault of 100 km² on the portion of an area of 10⁴ km² external to it, we find 353.2 events, while the same computation made including the negative contribution of the area internal to fault itself gives only 15.6 (less than 5%). The idea of an overall null balance is, then, substantially met: the phenomenon of clustering does not change the long term seismicity rate of a large area. This is in contrast with the opinion that trigger zones produce many more shocks than are missing from the shadows (Toda and Stein, 2003). We nevertheless consider Marsan (2003) who pointed out the practical difficulty of observing in nature the lack of events (the quiescence) mainly when looking at weakly active regions and short timescales.

5. Temporal behavior of the triggered seismicity

Let us now consider in more detail the dependence on time of the number of triggered events predicted by the present hypothesis. Kagan and Jackson (1991) and Dieterich (1994) found, in their empirical works, that shallow aftershocks, above a given magnitude cutoff, decay globally (for different distance intervals) in about 10 years and this time decay appears to systematically decrease with increasing depth. In our model the behavior in different zones depends on the different values of the stress changes, that decrease with the increasing distance from the edges of the fault. It can then be noted that the plateau of the plot of the seismicity rate *versus* time (see fig. 3) has a duration that increases with the increasing distance from the source. This result depends on our assumption of considering only aftershocks on the mainshock fault plane. It can also be biased by the fact that we consider only the effect of the static stress step $\Delta\tau$ caused by a mainshock and we neglect the interaction between other events. We also do not account for heterogeneities of frictional or constitutive prop-

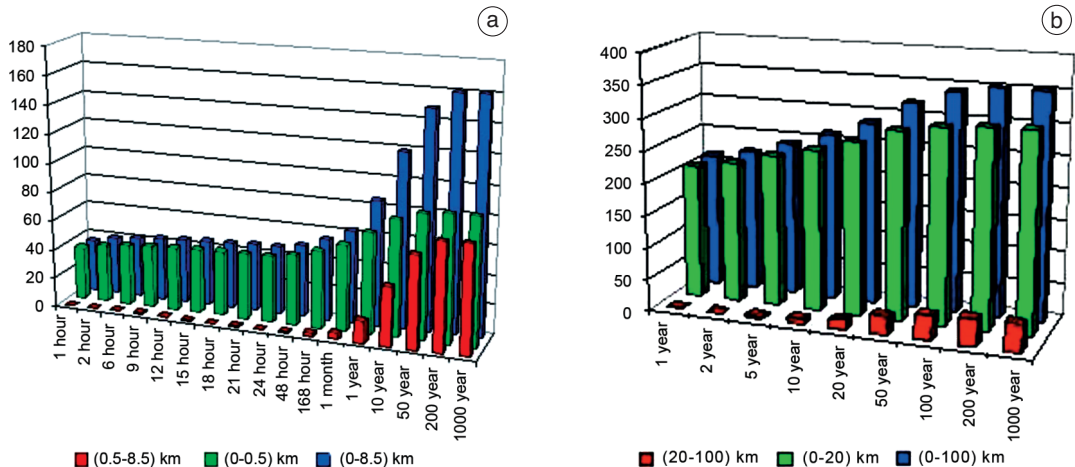


Fig. 5a,b. Representation of the total number of induced events (on vertical axis) in six different arbitrary slices (red, green and blue columns) around the principal fault after different intervals of time (on horizontal axis). The distances of the limits of the slices are computed from the perimeter of the source and the reference magnitude in the simulation is 6.0. Histogram a) represents the distribution nearest to the source at shorter intervals of time (the time scale is purely qualitative). Histogram b) represents the distribution farther from the fault at longer intervals of time.

erties nor for the dynamic stress concentration at the crack tip. In spite of that, we believe that our model can describe a behavior not so far from reality. Therefore we keep as our reference the model of faults given in table I again.

As shown in fig. 5a, in the close neighboring of the fault (within 0.5 km from the perimeter of the fault), most of the aftershocks occur within one day after the occurrence of the inducing earthquake, and more than 1/3 of them already occur in the first hour. This is clearly not the case for further distance from the source: in the slice between 0.5 and 8.5 km from the fault, about half of the activity is exploited within about one year. The further we go apart from the source, the longer the time interval necessary for the exploitation of the triggered activity. In fact, in the distant zone (from 20 to 100 km from the fault) most of the triggered events occur longer than 20 years after the inducing earthquake (fig. 5b). The behavior shown in fig. 5a,b is illustrated also by the snapshots of the spatial distribution of the seismicity rate at different times. Figures 6, 7 and 8 clearly show how the distribution of aftershocks resembles that of the stress change (fig. 2) and how the time decay is different in the different slices. These also show how shadow zones (where the seismicity is inhibited) have a longer effect in time. It should be noted again that we make the assumption that the aftershocks occur

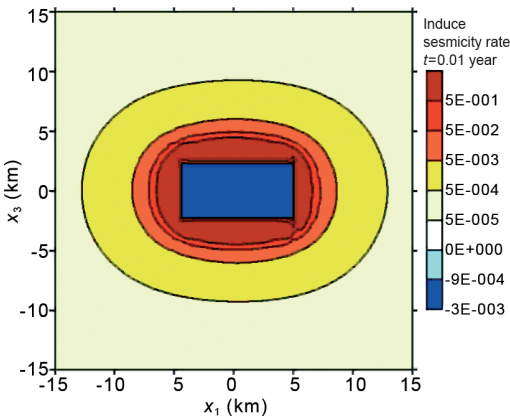


Fig. 6. Spatial distribution of the density of induced events at $t=0.01$ year.

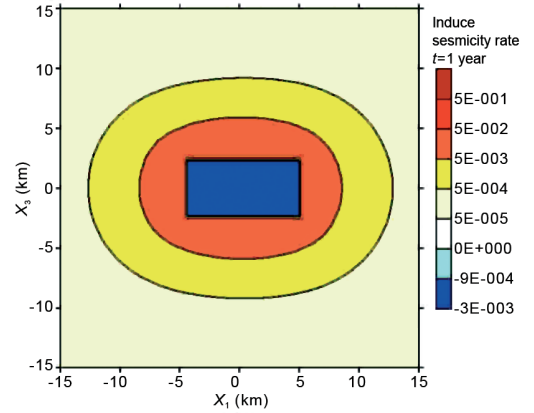


Fig. 7. Spatial distribution of the density of induced events at $t=1$ year.

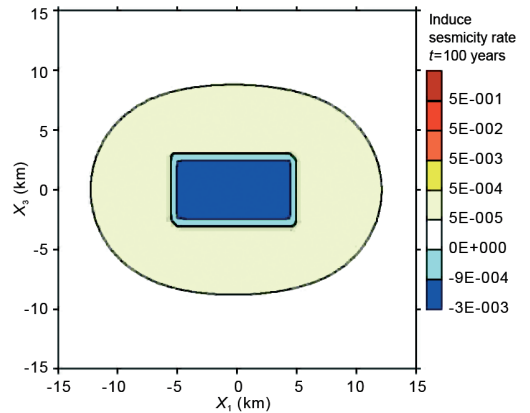


Fig. 8. Spatial distribution of the density of induced events at $t=100$ year.

mainly on the fault plane and that these have the same focal mechanism as the causative event.

We are now interested in checking if the generalized Omori Law is suitable to describe the temporal behavior of $R'(t)$ predicted by our model in every spatial slice and, in such case, by means of which free parameters. We refer to the generalized Omori Law as to the formula describing the decay of aftershock rate after a mainshock (Utsu *et al.*, 1995) through three free parameters a , c and p

Table VII. Results obtained for the bestfit of the sets of data reported in fig. 5a,b by the relation (5.2). The symbol rms indicates the standard deviation.

Slice	a [y^{p-1}]	c [y]	p	rms
(0÷0.5) km	3.24 ± 0.17	$0.37 \cdot 10^{-7} \pm 0.2 \cdot 10^{-6}$	1.02 ± 0.02	1.86
(0÷8.5) km	11.0 ± 0.76	$0.23 \cdot 10^{-7} \pm 0.2 \cdot 10^{-6}$	0.91 ± 0.032	9.16
(0.5÷8.5) km	83.7 ± 70.0	3.1 ± 1.96	1.44 ± 0.19	3.3
(0÷20) km	34.3 ± 8.0	$0.13 \cdot 10^{-1} \pm 0.019$	1.14 ± 0.077	9.22
(0÷100) km	36.8 ± 9.1	$0.15 \cdot 10^{-1} \pm 0.023$	1.12 ± 0.08	10.8
(20÷100) km	$12.5 \cdot 10^3 \pm 10.9 \cdot 10^3$	61.0 ± 5.43	$2.8 \pm 0.96 \cdot 10^{-3}$	0.45

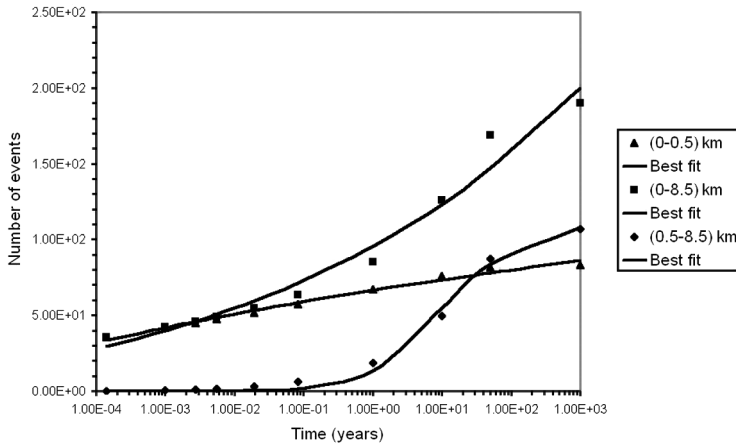


Fig. 9. Comparison between the temporal distribution of the cumulative number of induced events corresponding to synthetic data and the best fits obtained by the Omori Law for the closer slices from the source.

$$R'(t) = \frac{a}{(c+t)^p}. \quad (5.1)$$

For this purpose we make use of an algorithm for the least squares bestfit of the sets of data reported in fig. 5a,b by the integral of (5.1) over the time

$$y = \int_0^t R'(t') dt' = a \left[\frac{(c+t)^{1-p} - c^{1-p}}{1-p} \right]. \quad (5.2)$$

The results obtained for the bestfit are shown in table VII. It is evident that the values of the pa-

rameters obtained are quite different from case to case (for instance, the c parameter is extremely small for the slices closest to the fault edges), and that the standard deviation is larger for the slices that include the wider range of distances from the fault. These results show that the same set of parameters in eq. (5.2) can not fit the synthetic values obtained for different zones. Dieterich (1994) reached a similar conclusion for a source model represented by a circular shear crack. However, his analysis was limited to the spatial variation of p parameter letting $c=0$. The observation that the Omori p parameter varies spatially in real aftershocks sequences has been documented not only by

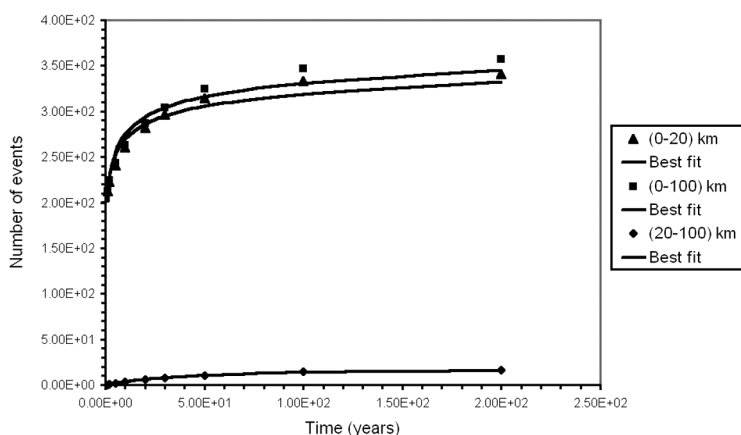


Fig. 10. Comparison between the temporal distribution of the cumulative number of induced events corresponding to synthetic data and the best fits obtained by the Omori Law for the farther slices from the source.

Dieterich (1994) but also in a number of papers (Wiemer and Katsumata, 1999; Wiemer, 2000; Enescu and Ito, 2002; Wiemer and Wyss, 2002; Wiemer *et al.*, 2002). Wiemer *et al.* (2002) reported a case of significantly different c values in the northern and southern Hector Mine aftershock volume, respectively. Enescu and Ito (2002) in their study of aftershock activity of the 2000 Western Tottory earthquake found c -values very close to 0 except for larger magnitude earthquakes, arguing that this feature could result from incompleteness of data, or might also reflect the complexity of the rupture process.

The situation is illustrated also by the plots of figs. 9 and 10, showing the graphical comparison between the synthetic data and the relative theoretical approximation for different slices.

6. Conclusions

In the development of this model some approximations have been made. Two of the most important are a) having neglected the dishomogeneity of the stress-drop on the fault, and consequently the very complicated pattern of slip (this must certainly influence the aftershock pattern: some aftershocks may even occur inside the fault, where the gradient of slip is

high); b) having neglected the interaction between subsequent events.

The first step of our modeling, concerning the stress change in the medium following an earthquake, has given results comparable to those of previous papers like, for instance, Chinnery (1961) and King and Cocco (2001). This substantial agreement, though the simplifications introduced in this specific case of rectangular source, supports the validity of our methodology.

Neglecting the interaction among seismic events and using the rate-and-state model of Dieterich (1994) entails that the total number of triggered events per unit area is proportional to the shear stress change. Nevertheless, the time constant by which the rate of events decreases has a large variability, ranging from a fraction of hour for a shear stress change of several MPa, up to tens of years for stress variations smaller than 1 MPa.

A first set of simulations carried out with a model of rectangular fault, based on the hypothesis that most of the aftershocks occur on the same plane as the fault, has led to the relationship between the main source parameters and the number of triggered events. The simulations lead to the conclusion that many earthquakes of small magnitude can produce, in total, a number of aftershocks comparable to that

of fewer larger earthquakes. Our model predicts, in fact, that, for a constant stress-drop, the number of triggered events is roughly proportional to the seismic moment at the power of $2/3$, or equivalent to 10^M , as observed on various occasions. Helmstetter (2003) and Helmstetter *et al.* (2005) suggest that this behavior can be related to the fractal structure of the spatial distribution of the seismicity and to the nature of earthquake interactions. For a constant stress-drop we also observe that the total number of triggered events is proportional to the area of the fault; on the other way, if the linear dimensions of the fault are kept constant, the number of triggered events is proportional to the stress-drop.

An important aspect of the model concerns the time-space behavior of the triggered seismicity. The simulations show that the application of the rate-state model and eq. (3.1) introduced by Dieterich (1994) to the static model of stress change has as a consequence something more complicate than the conventional Omori Law for the temporal decay of aftershock rate. In fact, the seismicity is expected to be intense (in accordance to the stress change) and have a maximum constant value, close to the fault edges of the primary earthquake, but the time necessary for the starting of decay is shorter. On the contrary, the seismicity expected at larger distances (up to 10 times the linear dimensions of the fault) is weak, but the time necessary for the decay from the maximum value is longer (even some tens of years). We can conclude that the temporal behavior of the aftershock rate varies in space. Moreover, the model predicts that the total long-term production of aftershocks, including wider areas, is not negligible respect to the short-term production. In this respect, the regional background seismicity could be interpreted as a sort of noise, or memory, due to the superposition of the effect of many older earthquakes. The time decay described by the popular Omori Law could be interpreted as an apparent average result of the contribution from the various areas on the plane containing the primary fault.

The rate-state model (Dieterich, 1994) used in our algorithm for the time distribution of the seismicity produced by a dislocation on a rec-

tangular fault, has allowed a partial justification of the Omori Law. We outline the following main conclusions obtained from the simulations:

- The seismic activity is very intense during the first few hours or days after the occurrence time of the primary earthquake, in the rectangular slice close to the edges of the fault.
- The immediate decay of this intense activity is followed by a period of time during which the activity at larger distance (of the order of the linear dimensions of the fault) is nearly constant. For longer time (several years) the total number of triggered events in the external zone becomes comparable with that of the most internal one.
- The number of events triggered over a long time scale at distances larger than the fault linear dimensions is significantly large.
- The seismic activity inside the fault drops to negligible values at the occurrence of the primary earthquake and returns to normal values only after a time much longer than the characteristic time (several tens of years in our simulations); for a long time range the number of events inhibited inside the fault is comparable to that of the events triggered outside: the net balance can be considered null; the seismicity rate changes in space and time, but does not influence the rate averaged over very large intervals of space and time.

In the light of the consequences derived from our model, both positive and negative perturbations of the seismic activities resulting from the interactions among earthquakes would last for a time in the order of decades. In this respect, it seems unrealistic to define and observe what is commonly called the «background seismicity» of an area.

We can conclude that, in spite of its very simple conception, the model here analyzed gives a physical justification both to the very popular phenomenon of short-term earthquake clustering (aftershocks and foreshocks in a strict sense) and to that of the long-term induction and quiescence, also observed in various occasions.

We are confident that the refinement of the model and its validations with real catalogs will lead to even more interesting consequences.

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